#### **Introduction to High Performance Computing**

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SDS406 – Fall semester, 2024 - 2025

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L09: Parallelizing PDEs with MPI,  $25^{th}$  November 2024

Partial Differential Equations (PDEs) arise naturally, e.g. describing time-evolution in physical systems

• Consider the PDE:

$$\partial_t u(t,\vec{r}) = \kappa \nabla_{\vec{r}}^2 u(t,\vec{r})$$

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• Consider the PDE:

$$\partial_t u(t, \vec{r}) = \kappa \nabla_{\vec{r}}^2 u(t, \vec{r})$$

- $u(t, \vec{r})$  can be a field of temperatures at position  $\vec{r}$  and time t. The above PDE describes the time-evolution of  $u(t, \vec{r})$ , given:
  - An initial condition,  $u(0, \vec{r})$
  - $\circ~$  A set of boundary conditions, e.g.:
    - $u(t, \vec{r}) = 0$  for  $r^2 > R^2$  (zero if outside of a sphere with radius R);
    - $u(t, \vec{r} + \hat{e}_i L) = u(t, \vec{r})$  for i = 1, 2, 3 (periodic within box of side L).

Numerical solution of PDEs requires first writing the PDE in discrete form

• We will restrict to the 2-dimensional PDE (2 spatial dimensions, x and y, plus time, t):

$$\partial_t \mathfrak{u}(t,x,y) = \kappa[\partial_x^2 \mathfrak{u}(t,x,y) + \partial_y^2 \mathfrak{u}(t,x,y)]$$

• One choice of discretizing the derivative

$$f(x+\delta) = f(x) + \frac{\partial f(x)}{\partial x}\delta + \frac{\partial^2 f(x)}{\partial x^2}\frac{\delta^2}{2} + O(\delta^3) \Rightarrow$$
$$\frac{\partial f(x)}{\partial x} = \frac{f(x+\delta) - f(x)}{\delta} + O(\delta)$$

• For second derivative, take the central difference:

$$f(x+\delta) = f(x) + \frac{\partial f(x)}{\partial x}\delta + \frac{\partial^2 f(x)}{\partial x^2}\frac{\delta^2}{2} + O(\delta^3)$$
$$f(x-\delta) = f(x) - \frac{\partial f(x)}{\partial x}\delta + \frac{\partial^2 f(x)}{\partial x^2}\frac{\delta^2}{2} + O(\delta^3)$$

• Add both sides:

$$\frac{\partial^2 f(x)}{\partial x^2} \simeq \frac{f(x+\delta) - 2f(x) + f(x-\delta)}{\delta^2}$$

• Back to the heat equation:

$$\begin{aligned} \partial_{t} u(t,x,y) &= \kappa [\partial_{x}^{2} u(t,x,y) + \partial_{y}^{2} u(t,x,y)] \Rightarrow \\ & \frac{u(t+\tau,x,y) - u(t,x,y)}{\tau} = \\ \frac{\kappa}{2} [u(t,x+h,y) + u(t,x-h,y) + u(t,x,y+h) + u(t,x,y-h) - 4u(t,x,y)] \end{aligned}$$

This allows us to write the heat equation iteration:

 $u(t+\tau, x, y) = u(t, x, y)(1-4\alpha) + \alpha H[u(t, x, y)]$ 

where  $\alpha = \frac{\tau \kappa}{h^2}$  and H, the hopping term:

 $H[\mathfrak{u}(t,x,y)] = \mathfrak{u}(t,x+h,y) + \mathfrak{u}(t,x-h,y) + \mathfrak{u}(t,x,y+h) + \mathfrak{u}(t,x,y-h).$ 

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H[u(t,x,y)] = u(t,x+h,y) + u(t,x-h,y) + u(t,x,y+h) + u(t,x,y-h).

• Note that stability analysis requires:

 $\alpha < 0.25$ 

for stability of the iterations.

#### Serial code for heat equation

• Copy the directory:

/onyx/data/sds406f24/l09/ex01/

to your home directory.

- Included is a C source code file and a python script
- Study the C file, complete the missing parts, and run it using:

[user@front02 l09]\$ srun -n 1 -p p100 ./heat 256 6000 0.2

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- Arguments are:
  - L, the spatial extent,
  - T, the number of time iterations, and
  - alpha, defined as in the previous slide

#### Serial code for heat equation; Some implementation details

#### Indexing and 2-dimensional arrays

• The 2-dimensional array that holds u(x, y) is stored as 1-dimensional array in memory, e.g.:  $v[M][N] \rightarrow v[M*N]$ 

#### Serial code for heat equation; Some implementation details

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- We have decided the order of the indices to be such that *x* runs fastest:

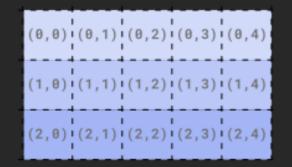
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#### Serial code for heat equation; Some implementation details

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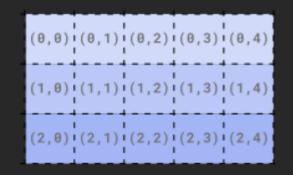
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• The C macro IDX() saves us some typing:

#define IDX(y, x) ((L+(y))%L)\*L + ((L+(x))%L)

Not a function! Macros are expanded at compile time

(0,0) (0,1) (0,2) (0,3) (0,4) (1,0) (1,1) (1,2) (1,3) (1,4) (2,0) (2,1) (2,2) (2,3) (2,4)

#### Serial code for heat equation; Some implementation details

#### Boundary and initial conditions

- 1. All points with x = 0 and x = L-1 are set to 1
- 2. All points on y = L/2 are set to 1
- 3. All other points are zero
- 4. Conditions 1. and 2. are enforced in every iteration

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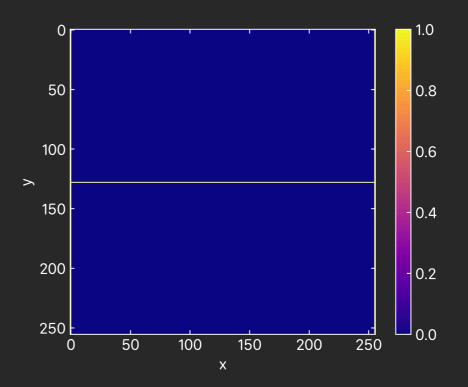
```
/***
 * Set the boundary condition for v[L*L]
 ***/
void
boundary_condition(double *v)
{
 for(int y=0; y<L; y++) {
   v[IDX(y, 0)] = 1;
   v[IDX(y, L-1)] = 1;
  }
 for(int x=0; x<L; x++) {
   v[IDX(L/2, x)] = 1;
  }
 return;
}
....
for(int i=0; i<T; i++) {
   /* Apply one time-step */
   update(v0, alpha);
   /* Update boundaries */
   boundary_condition(v0);
 }</pre>
```

#### Serial code for heat equation; Some implementation details

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- The final distribution is stored in v0.txt
- You can plot v0.txt using the python file plot.py

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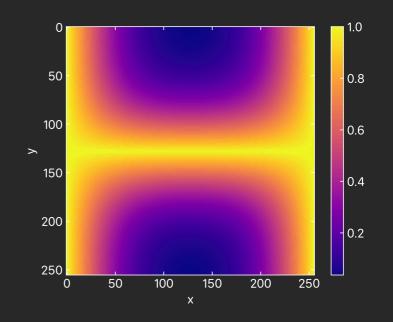
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# **Parallelization using MPI**

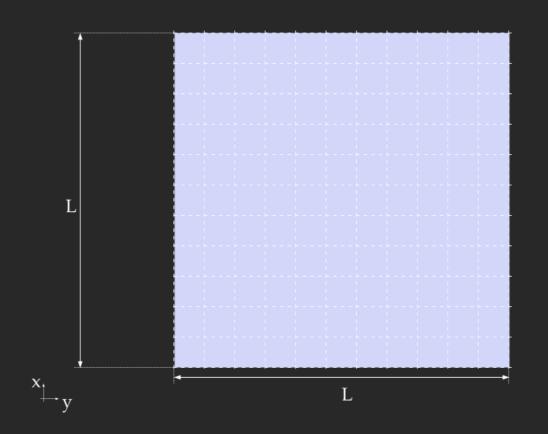
#### **Domain Decomposition of Partial Differential Equations (PDEs) using MPI**

- Parallelizing the 2-dimensional heat equation with MPI
  - Parallelizing the slow-running coordinate
  - Domain decomposition
  - Boundary exchange
  - Boundary conditions

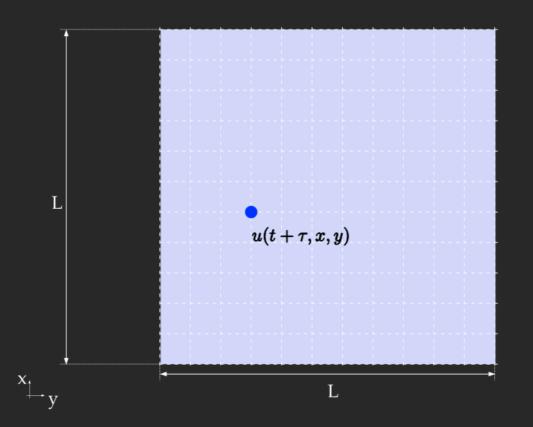
- Domain decomposition is referred to when we parallelize a problem by distributing the domain over processes
- In the case of the heat equation, we begin with an initial condition u(t = 0, x, y) and every iteration depends on the previous time ⇒ cannot parallelize in time coordinate
- However we *can* parallelize over the spatial domain

 $u(t+\tau, x, y) = u(t, x, y)(1-4\alpha) + \alpha H[u(t, x, y)],$ 

 $u(t+\tau, x, y) = u(t, x, y)(1-4\alpha) + \alpha H[u(t, x, y)],$ 

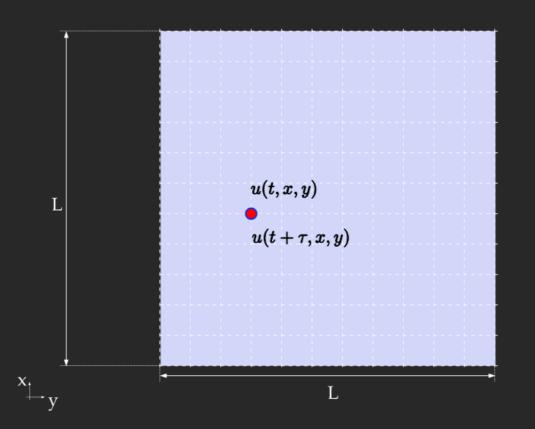


 $u(t+\tau, x, y) = u(t, x, y)(1-4\alpha) + \alpha H[u(t, x, y)],$ 



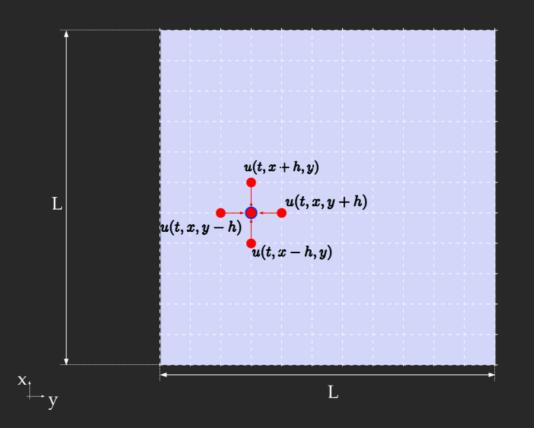
• Note the dependencies of each grid-point

 $u(t+\tau,x,y) = u(t,x,y)(1-4\alpha) + \alpha H[u(t,x,y)],$ 



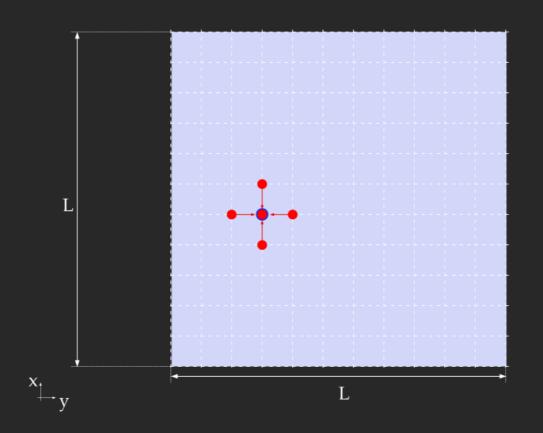
• Each grid-point (x, y), for  $t + \tau$ , requires (x, y) at t

 $u(t+\tau,x,y) = u(t,x,y)(1-4\alpha) + \alpha H[u(t,x,y)],$ 

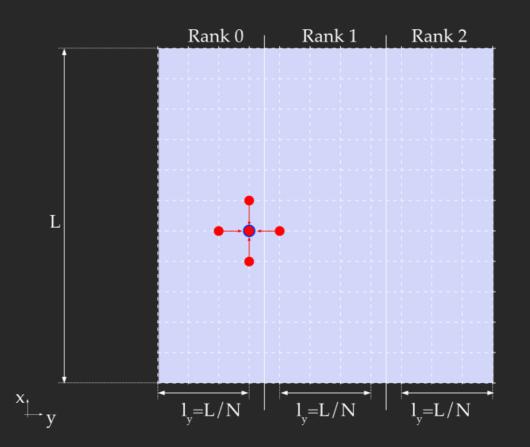


• ...as well as (x, y)'s four nearest neighbors

 $u(t+\tau, x, y) = u(t, x, y)(1-4\alpha) + \alpha H[u(t, x, y)],$ 

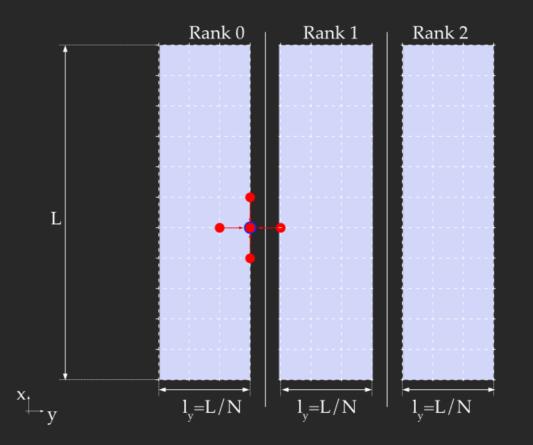


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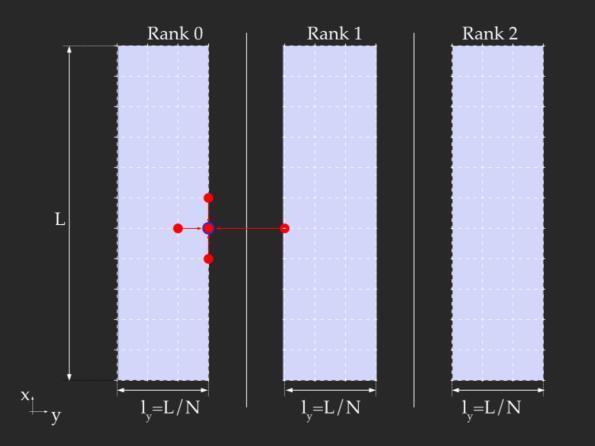
• Assume we are partitioning over the y direction

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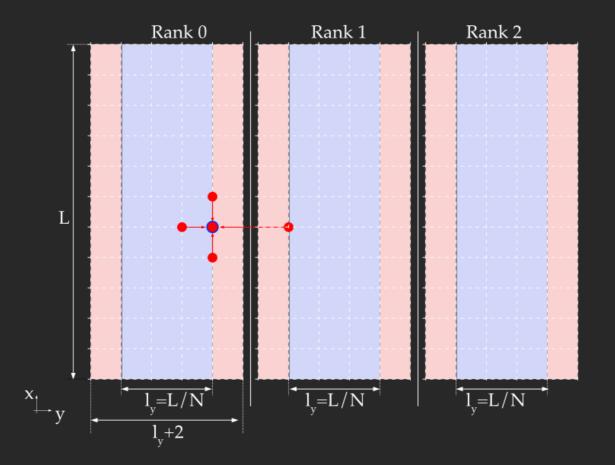
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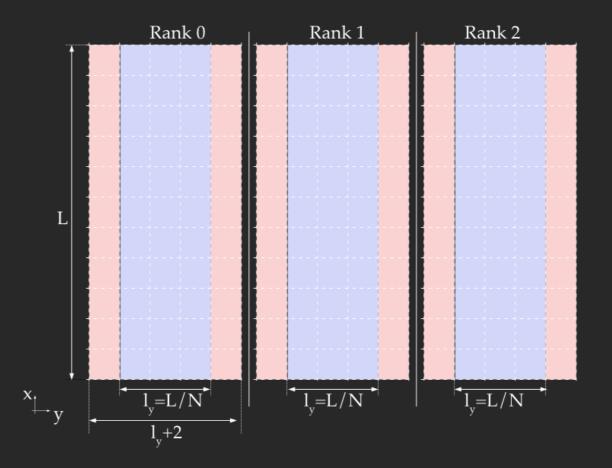
- Assume we are partitioning over the y direction
- For all processes, there are missing dependencies for y=0 and  $y=l_y-1$

 $u(t+\tau,x,y) = u(t,x,y)(1-4\alpha) + \alpha H[u(t,x,y)],$ 



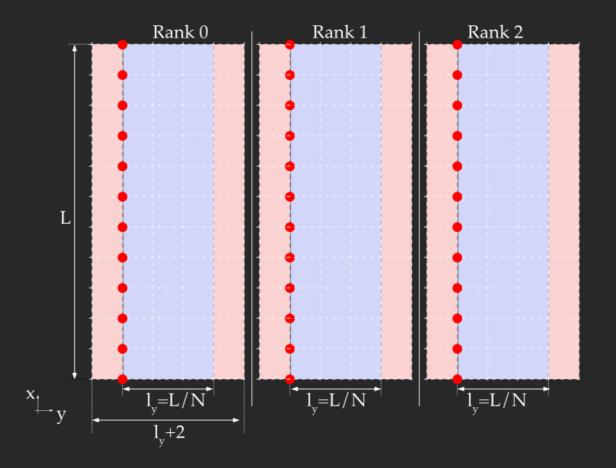
• The solution is to allocate "halos". Extra space to store the neighboring rank's data

 $u(t+\tau,x,y) = u(t,x,y)(1-4\alpha) + \alpha H[u(t,x,y)],$ 



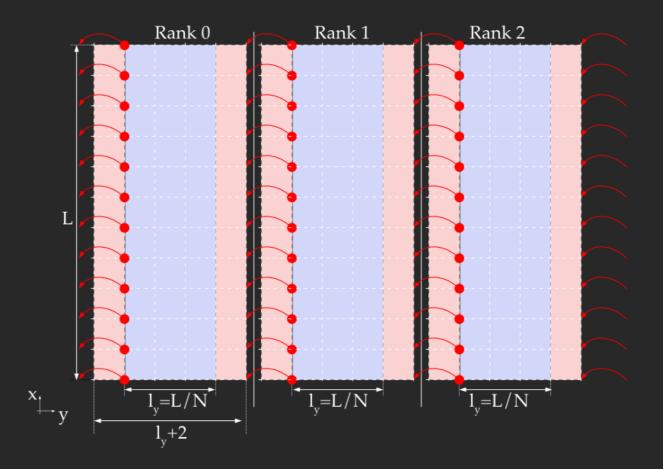
• Communication proceeds in two steps

 $u(t+\tau,x,y) = u(t,x,y)(1-4\alpha) + \alpha H[u(t,x,y)],$ 



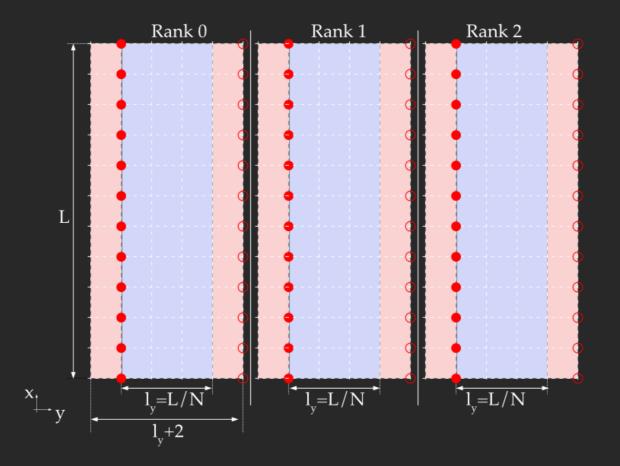
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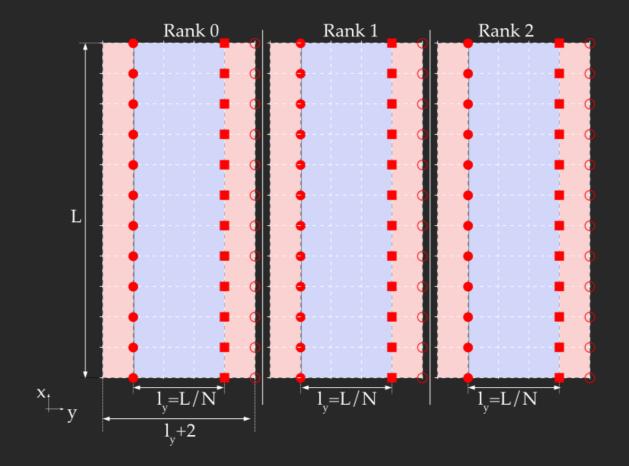
• The "backwards boundaries" are communicated "backwards" and received from the "forwards" process

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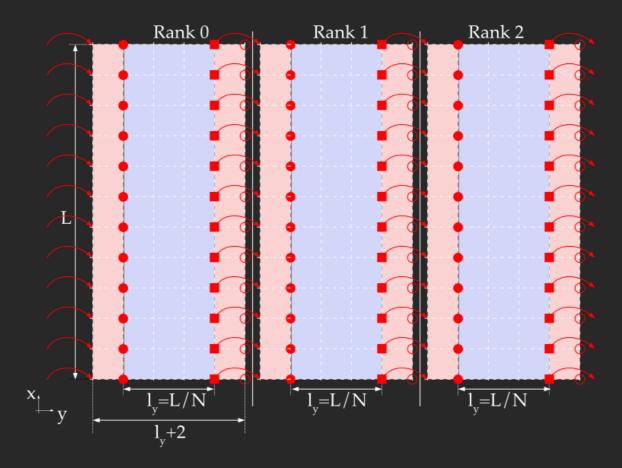


• Now the halos include the (rank+1)'s boundaries; the neighbors for  $y = l_y$  are available

 $\overline{u(t+\tau,x,y)} = u(t,x,y)(1-4\alpha) + \alpha H[u(t,x,y)],$ 

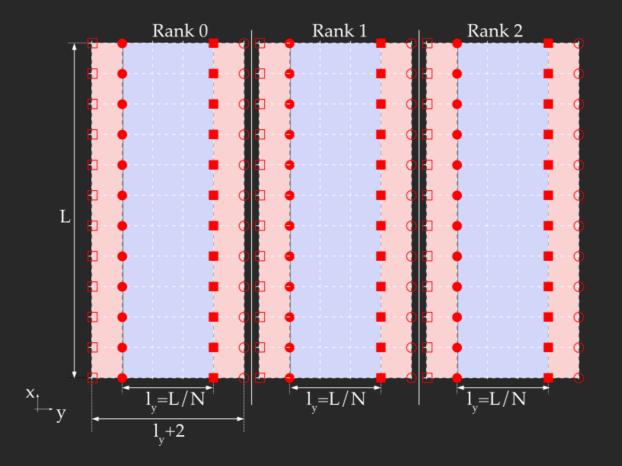


 $\mathfrak{u}(t+\tau,x,y) = \mathfrak{u}(t,x,y)(1-4\alpha) + \alpha H[\mathfrak{u}(t,x,y)],$ 



• The "forward boundaries" are communicated "forwards" and received from the "backwards" process

 $u(t+\tau, x, y) = u(t, x, y)(1-4\alpha) + \alpha H[u(t, x, y)],$ 



• Now the halos include the (rank-1)'s boundaries; the neighbors for y = -1 are available

#### MPI parallelization of the heat equation

• Copy the directory:

/onyx/data/sds406f24/l09/ex02/

to your home directory.

- Included is a C file that includes MPI
- It is incomplete, and you must complete the missing sections
- Once done you can run it using the heat.sh Slurm script
- It should produce exactly the same result as the serial version

#### MPI parallelization of the heat equation

1st todo

- Fix the IDX() macro
- Remember, the extent in x is the same
- The extent in y is now shorter. But also, there is space allocated before y=0 and after y=ly-1

#### MPI parallelization of the heat equation

#### **2nd TODO**

- Fix the write-function
- We want each rank to write, in order, its part of the array, appending to a single file
- You need to ensure serialization of the rank writing

# **Domain decomposition for the heat equation** MPI parallelization of the heat equation

#### 3rd todo

- Fix the function that sets the boundary conditions
- Use the same boundary conditions as in the serial version
- You need to think about which rank holds which grid-points

#### MPI parallelization of the heat equation

#### 4th Todo

- Fix the function that communicates the boundaries, filling in the "halos" of each process
- You need to do two communications
  - A Send to the next neighbor/receive from the previous neighbor
  - $\circ~$  A Send to the previous neighbor/receive from the next neighbor

#### **Domain decomposition for the heat equation** MPI parallelization of the heat equation

#### 5th TODO

- Complete the update of the field v in function update()
- How will this be different from the serial version?

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#### When done

- Compile with mpicc
- Run with heat.sh
- The resulting output v0.txt should be *identical* to the serial version

#### MPI parallelization of the heat equation

#### 5th TODO

- Complete the update of the field v in function update()
- How will this be different from the serial version?

#### When done

- Compile with mpicc
- Run with heat.sh
- The resulting output v0.txt should be identical to the serial version
- E.g.

[ikoutsou@front02 ex02]\$ diff v0.txt ../ex01/v0.txt
[ikoutsou@front02 ex02]\$

#### No output means no differences found

#### Strong scaling of the MPI-parallelized heat equation solution

For next lesson:

 $\Rightarrow$  Plot against inverse time-to-solution rather than speed-up

 $\Rightarrow$  Include in the plot the scalar version of  ${\tt ex01}$  at  $n_{proc}=1$